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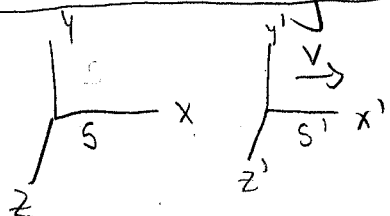
Review ; Postulates of SR:

- ① The speed of light is constant in all reference frames
- ② The laws of physics are the same in all inertial frames.

3 Fundamental consequences

- ① Loss of Simultaneity
- ★ ② Time Dilation:
- ★ ③ Length Contraction

Described by Lorentz Transformations :



$$\left. \begin{aligned} x' &= \gamma(x - vt) \\ y' &= y \\ z' &= z \\ t' &= \gamma(t - vx/c^2) \end{aligned} \right\}$$

Consider
Inverse
Transformations

$$\gamma = \frac{1}{\sqrt{1 - v^2/c^2}} \geq 1$$

Time Dilation

Two Events happen at same position in S
for simplicity, with time interval τ_0
 $x=0$ for simplicity, with time interval τ_0
 $t_1=0$ $t_2=\tau_0$

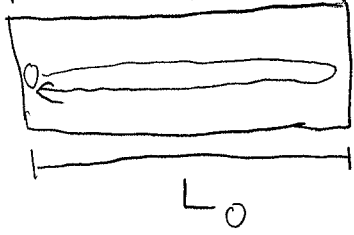
In S' , from LT : $\tau' = \gamma \tau_0$

Time interval is longer, clock runs "slow".

Looking from other frame not a contradiction. Result relies on clocks being at rest in a given frame.
- More intuitive derivation on HW. "Light clock"

Length Contraction

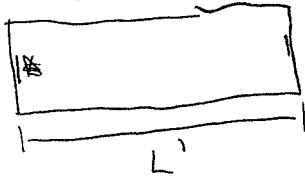
Measure a round trip of light pulse on train
In rest frame of train



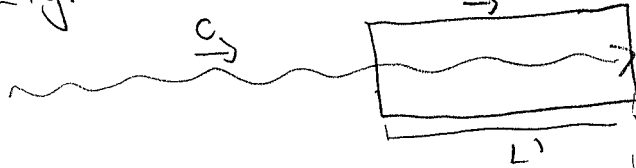
$$\tau_0 = \frac{2L_0}{c} \Rightarrow L_0 = \frac{c\tau_0}{2}$$

In Ground Frame

① Light Emitted



② Light hits back @ Δt_1

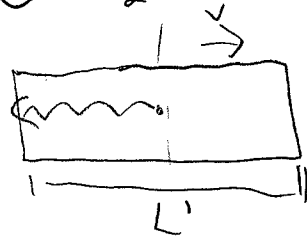


$$\Delta t_1 = \frac{L'}{c} + \frac{\Delta t_1 v}{c}$$

$$\Delta t_1 (1 - v/c) = \frac{L'}{c}$$

$$\Delta t_1 = \frac{L'}{c-v}$$

③ Light hits front @ Δt_2



$$\Delta t_2 = \frac{L'}{c} - \frac{\Delta t_2 v}{c}$$

$$\Delta t_2 = \frac{L'}{c+v}$$

Total Round Trip: $\tau_G = \Delta t_1 + \Delta t_2 = \frac{L'}{c-v} + \frac{L'}{c+v}$

$$\tau_G = \frac{2L'c}{c^2 - v^2} = \frac{2L'}{c} \gamma^2$$

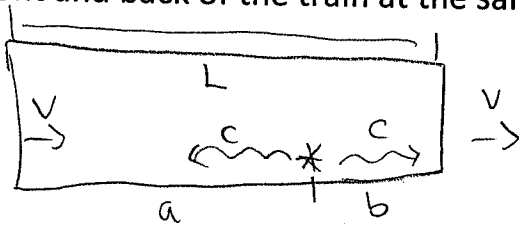
But the Ground observer sees the train clock running slower

$$\tau_G = \gamma \tau_0 \Rightarrow \frac{2L'}{c} \gamma^2 = \gamma \left(\frac{2L_0}{c} \right)$$

$$\boxed{L' = L_0 / \gamma}$$

← Length shortens as seen from moving frame.

A train with proper length L moves to at speed v . Where should the light source be located on the train, so than an observer *on the ground* sees the photons hit the front and back of the train at the same time?



$L = a + b$ (Train Frame)
 $a' = a/\gamma, b' = b/\gamma$ (Length Contraction)

Ground Frame | $t'_b = \frac{b/\gamma}{c} + \frac{t_b v}{c}$

$t'_b = \frac{b/\gamma}{c} \frac{1}{(1-v/c)} = \frac{b}{\gamma(c-v)}$

$t_a = \frac{a/\gamma}{c} - \frac{t_a v}{c} \Rightarrow t_a = \frac{a}{\gamma c} \frac{1}{(1+v/c)} = \frac{a}{\gamma(c+v)}$

Same Arrival Time in Ground Frame

$t'_a = t'_b \quad \frac{a}{c+v} = \frac{b}{c-v}$

Use: $L = a + b \Rightarrow$

$\frac{L-b}{c+v} = \frac{b}{c-v}$
 $a = L - \frac{L(c-v)}{2c} = \frac{2cL - cL + Lv}{2c} = \frac{L(c+v)}{2c} = a$
 $b = \frac{L(c-v)}{2c}$

Feedback: Today's lecture is

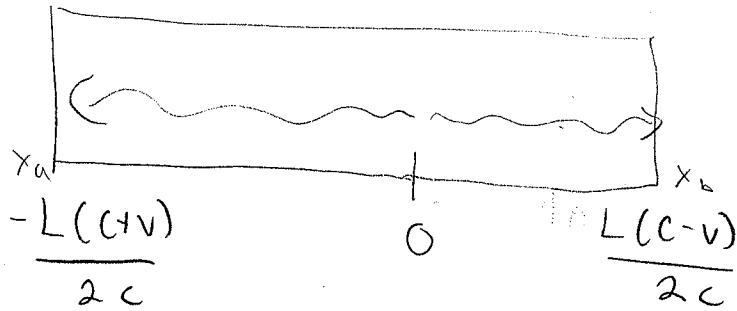
Too fast

about right

too slow

Other comments: _____

Exercise! Use results to confirm L-contraction using Lorentz transforms



$$t'_a = t'_b$$

$$-\left(\frac{L(c+v)}{2c}\right) = \gamma(x'_a - vt'_a)$$

$$\frac{L(c-v)}{2c} = \gamma(x'_b - vt'_b)$$

Diff! $\frac{L(c-v)}{2c} + \frac{L(c+v)}{2c} = \gamma(x'_b - x'_a)$

$$L \frac{2c}{2c} = \gamma L' \Rightarrow \boxed{L' = L/\gamma}$$

Real Life Examples:

Muons: Created in upper atmosphere, a live w/ lifetime 2 ms. with $v = (1 - 2 \times 10^{-5})c$. If they are created 50 km up, do they hit the ground.

Newton would Have said No: $d = vT \approx 600 \text{ m} < 50 \text{ km}$.

SR from Observer's Perspective: $d = v(\gamma T) = \frac{600 \text{ m} \gamma}{\sqrt{1 - (1 - 2 \times 10^{-5})^2}} \approx \frac{600 \text{ m}}{160} \approx 100 \text{ km}$

Time Dilation \rightarrow Yes!

SR from Muon's Perspective: $d/\gamma = vT \Rightarrow T_E = \frac{d}{\gamma v} = \frac{50 \text{ km}}{160 (3 \times 10^8)} \approx 1 \text{ ms}$

Length Contraction Yes! \rightarrow

World's Best Atomic Clock:

Measure a tick frequency

$$\nu_0 = \frac{1}{\tau_0}$$

→ rest frame

$$\nu' = \frac{1}{\tau'} = \frac{1}{\gamma \tau_0} = \frac{\nu_0}{\gamma}$$

→ moving frame

How stable does a slow clock need to be?
Measure a relative shift in frequency

$$\frac{\Delta \nu}{\nu_0} = \frac{\nu' - \nu_0}{\nu_0} = \frac{\nu_0 (1 - \frac{1}{\gamma})}{\nu_0} = 1 - \sqrt{1 - v^2/c^2}$$

$$\approx 1 - (1 - \frac{1}{2} v^2/c^2)$$

$$\boxed{\frac{\Delta \nu}{\nu_0} \approx \frac{1}{2} \frac{v^2}{c^2}}$$

for 1 m/s $\frac{\Delta \nu}{\nu_0} \approx 0.5 \times \frac{1}{9 \times 10^{16}} \approx 5 \times 10^{-18}$

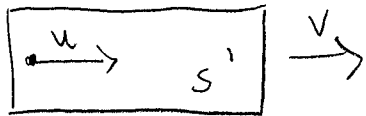
Just at the limit of world's most precise atomic clocks.
(Lose 1 sec is 4 billion years)
Chou et al (Science 2010)

measure $v = 5$ m/s $\frac{\Delta \nu}{\nu} \approx 1.4 \times 10^{-16}$

Phys 379: Lecture 9

- ① Real Life Examples from last lecture (Muons + clocks)
- ② Velocity Addition
- ③ Doppler Effect

Velocity Addition



What speed does the ball have relative to the ground?



EX: $u = \frac{1}{2}c$, $v = \frac{3}{4}c$

Newton would say: $v_g = u + v = \frac{5}{4}c > c$ X

Derive Relativistic Result from Lorentz Transformations

$$u = \frac{\Delta x'}{\Delta t'}$$

and

$$v_g = \frac{\Delta x}{\Delta t}$$

$$\Delta x = \gamma(\Delta x' + v\Delta t')$$

$$\Delta t = \gamma(\Delta t' + v\Delta x'/c^2)$$

Then:
$$v_g = \frac{\Delta x' + v\Delta t'}{\Delta t' + v\Delta x'/c^2} = \frac{\overbrace{\Delta x'/\Delta t'}^u + v}{1 + \underbrace{v/c^2}_{u} \Delta x'/\Delta t'}$$

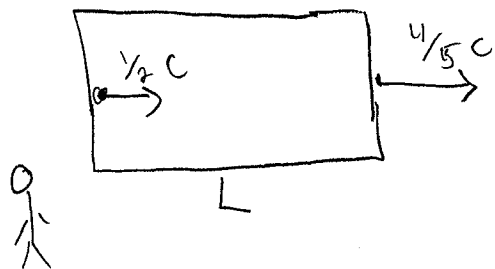
$$v_g = \frac{u + v}{1 + uv/c^2}$$

Important Properties:

- ✓ ① symmetric ($u \Leftrightarrow v$)
- ✓ ② $uv \ll c^2$, $v_g \approx u + v$ (Galilean velocity Transform)
- ✓ ③ If $u, v = c$, $v_g = c$ (speed of light same in all frames)

EX from Above:
$$v_g = \frac{(\frac{1}{2} + \frac{3}{4})c}{1 + (\frac{1}{2})(\frac{3}{4})c^2/c^2} = \frac{5/4c}{11/8} = \frac{10}{11}c$$

A ball is launched with speed $u = \frac{1}{2}c$ from the back of a train car with proper length L . If the train travels with speed $v = \frac{4}{5}c$, after how much time, according to an observer on the ground, does it take for the ball to hit the front of the train car? How far does it travel according to this observer?



Velocity of Ball relative to observer:

$$V_b = \frac{\frac{1}{2}c + \frac{4}{5}c}{1 + (\frac{1}{2})(\frac{4}{5})\frac{c^2}{c^2}} = \frac{\frac{13}{10}c}{\frac{7}{5}} = \frac{13}{14}c$$

The ball travels the contracted length of the train plus the extra distance it travels.

$$\gamma = \frac{1}{\sqrt{1 - \frac{16}{25}}} = \frac{5}{3}$$

$$t_b = \frac{L/\gamma}{V_b} + \frac{t_b v}{V_b} \Rightarrow t_b = \frac{L}{\gamma} \frac{1}{V_b - v} = L \left(\frac{3}{5}\right) \frac{1}{\left(\frac{13}{14} - \frac{4}{5}\right)c}$$

$$= L \left(\frac{3}{5}\right) \frac{70}{93} = \boxed{\frac{14}{3} L/c}$$

The distance is simply

$$\Delta x_b = V_b t_b = \left(\frac{13}{14}c\right) \left(\frac{14}{3} \frac{L}{c}\right) = \boxed{\frac{13}{3} L}$$

Same Procedure Gives Transverse Velocities

$$v_y = \frac{\Delta y}{\Delta t} = \frac{\Delta y'}{\gamma(\Delta t' + v \Delta x'/c^2)}$$

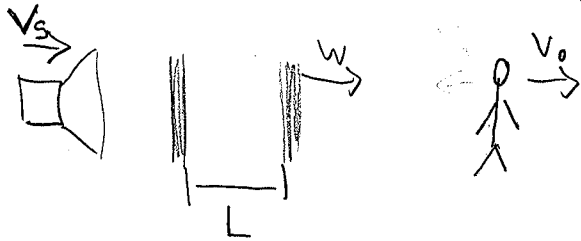
$$v_y = \frac{u_y'}{\gamma(1 + v u_x/c^2)}$$

Limerick: For a bullet, a train, and a gun
 Adding the speeds can be fun
 Take a trip down the path
 Paved with Einstein's new math,
 Where a half plus a half isn't one.

Exercise: Ball on Train

Doppler Effect:

In sound (Matter):



Period

$$T_o = \frac{1}{f_o}$$

Frequency

Wavelength

$$\lambda_o = \frac{w}{f_o} \leftarrow \text{speed}$$

Moving source: Distance between pulses compressed (wavelength modified)

$$\lambda_o = \lambda_o - v T = w/v_o - v/v_o = \frac{1}{v_o} (w - v)$$

$$v_o = v_o \frac{1}{1 - v_s/w} \approx v_o (1 + v_s/w)$$

Moving observer: Relative velocity of wave changes

$$v_o = \frac{(w - v_o)}{\lambda_o} = v_o (1 - v_o/w)$$

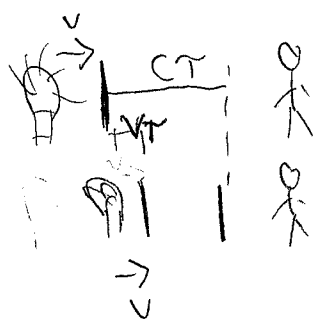
- ① Doppler Effect (classical + Relativistic)
- ② SR Paradoxes

Classical Doppler Effect

$$v_D = v_0 \left(\frac{1 - v_o/w}{1 + v_s/w} \right)$$

Relativistic Doppler Effect

Consider flashing light source + observer.
with relative velocity v



Frequency Modified Due to

① Time Dilation

$$\tau = \gamma \tau_0$$

② Wave length Compression

$$\lambda_D = c\tau - v\tau = (c - v)\tau$$

$$v_D = \frac{c}{\lambda_D} = \frac{c}{c - v} \frac{1}{\gamma \tau_0} = \frac{1}{1 - v/c} \frac{\sqrt{1 - v^2/c^2}}{\tau_0} = \frac{\sqrt{(1 + v/c)(1 - v/c)}}{(1 - v/c)} \frac{1}{\tau_0}$$

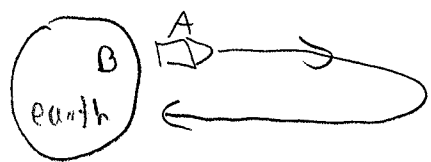
$$v_D = v_0 \sqrt{\frac{(1 + v/c)}{(1 - v/c)}}$$

Checks ; ① If $v > 0$, $v_D/v_0 > 1$. Light is "blueshifted"

② $v/c \ll 1 \Rightarrow v_D \approx v_0 (1 + \frac{1}{2}v/c)(1 + \frac{1}{2}v/c) \approx v_0 (1 + v/c)$

Agrees w/ Classical result

SR Paradoxes : Who is Younger?



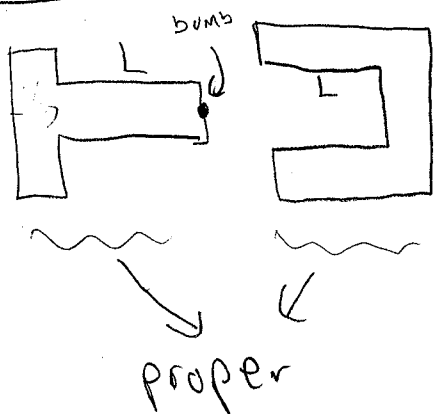
A sees B's clock run slow, + vice-versa.

But the situation is not symmetric. A must accelerate to turn around.

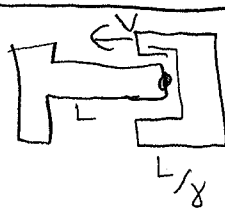
A is younger

While A is accelerating, sees A's clock runs extra slow from GR red shift time dilation.

T and U blocks : Travel toward each other with relative velocity v .



In T-frame:



U is contracted
Bomb explodes!

In U-frame



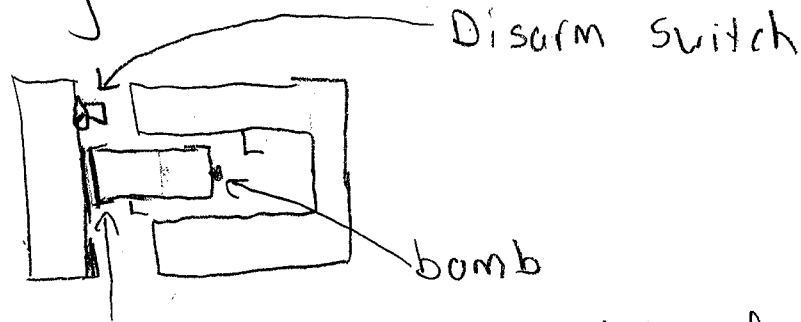
will it be stopped?

But signal has to propagate to front, to stop bomb. At best info travels at c .
 Bomb will explode if $\frac{L - L/\gamma}{v} < \frac{L}{c}$ ← time it takes stopping signal to propagate down block

← time it takes bomb to go remaining distance
 $1 - \frac{1}{\gamma} < \beta \Rightarrow 1 - \sqrt{1 - \beta^2} < \beta \Rightarrow 1 - \beta < \frac{\sqrt{1 - \beta^2}}{\sqrt{(1 - \beta)(1 + \beta)}}$
 $\Rightarrow \sqrt{1 - \beta} < \sqrt{1 + \beta} \checkmark$ True for any β !

The way we have formulated the problem implies the "T-block" will stretch until the bomb makes contact.

If you find this hard to accept, consider the following variation of the T-block,



This part can slide forward if the back stops.

Now the condition of the bomb exploding depends on whether or not the disarm switch hits first. But in the "U-block" frame, our argument from above still holds. The signal from the disarm switch can travel to the bomb at a ^{max} velocity of c . It is still true that this is not fast enough to catch up with the bomb.

So whether it is delayed disarm signals or stretching molecular bonds, the resolution to the paradox is the same. All observers agree that the bomb explodes.